

## Problem 1)

a) Draw a thin pillbox at the interface, with the top facet of the pill-box alone, and its bottom facet below, the interface. The surface integral corresponding to  $\vec{D} \cdot \vec{D}$  is  $\oint \vec{D} \cdot d\vec{s} = [\vec{D}_1(r^+, t) - \vec{D}_1(r^-, t)] \Delta s$ , with  $\Delta s$  being the surface area of the top (or bottom) facets. The total charge contained within the pill-box is  $J_{s-free}(r^+, t) \Delta s$ . Setting these equal to each other, we find:  $\vec{D}_1(r^+, t) - \vec{D}_1(r^-, t) = J_{s-free}(r^+, t)$ .

b) Draw a thin, rectangular loop at the interface. The short legs of the loop are  $\perp$  to the interface. The long legs are parallel to the interface, one being above, the other one below. The integral of  $\partial \vec{D} / \partial t$  over the area of the loop may be ignored (because the loop is arbitrarily thin). We distinguish two cases:

Case I : Area of the loop is parallel to  $\vec{J}_{s-free}(r^+, t)$  at the location of interest.

In this case, the surface current makes no contribution to the integral over the loop surface. We then have  $\oint \vec{H} \cdot d\vec{l} = [\vec{H}_1(r^+, t) - \vec{H}_1(r^-, t)] \cdot \vec{\Delta l} = 0$   
 $\Rightarrow \vec{H}_1(r^+, t) = \vec{H}_1(r^-, t)$ . In words, the component of the  $\vec{H}$  field that is parallel to the interface AND also parallel to  $\vec{J}_{s-free}(r^+, t)$  is continuous across the interface.

Case II : area of the loop is  $\perp$  to the direction of  $\vec{J}_{s-free}(r^+, t)$ . The contribution of the surface current to the integral over the loop area is  $J_{s-free} \Delta l$ . Setting this equal to  $\oint \vec{H} \cdot d\vec{l}$  we find:  $\vec{H}_1(r^+, t) - \vec{H}_1(r^-, t) = J_{s-free} \vec{\Delta l}$ . In words, the component of  $\vec{H}$  that is parallel to the interface AND perpendicular to  $\vec{J}_{s-free}$  has a discontinuity equal to  $J_{s-free}$  at the interface.

c) Same method as part (b) yields  $\vec{E}_1(r^+, t) = \vec{E}_1(r^-, t)$ .

d) Same method as part (a) yields  $\vec{B}_1(r^+, t) = \vec{B}_1(r^-, t)$ .

Problem 2)

a)  $\vec{H}(x, t) = H_0 \cos\{\omega[t - n(\omega)x/c]\} \hat{z}$

Maxwell's Third equation  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \Rightarrow \frac{\partial E_y}{\partial x} \hat{z} = -\mu_0 \frac{\partial H_z}{\partial t} \hat{z} \Rightarrow$

$$E_0 \cos n(\omega)t/c \sin\{\omega[t - n(\omega)x/c]\} = \mu_0 H_0 \omega \sin\{\omega[t - n(\omega)x/c]\} \Rightarrow$$

$$\underbrace{H_0}_{= \frac{n(\omega)}{\mu_0 c} E_0} = \frac{n(\omega)}{Z_0} E_0 \quad \checkmark$$

b)  $\vec{S}(x, t) = \vec{E}(x, t) \times \vec{H}(x, t) = E_0^2 H_0 \cos^2\{\omega[t - n(\omega)x/c]\} \hat{x} \quad \checkmark$

$$\langle \vec{S}(x, t) \rangle = E_0^2 H_0 \langle \cos^2\{\omega[t - n(\omega)x/c]\} \rangle \hat{x} = \underbrace{\frac{1}{2} \frac{n(\omega)}{Z_0} E_0^2 \hat{x}}_{\checkmark}$$

c)  $\vec{E}_1(x, t) + \vec{E}_2(x, t) = E_0 \hat{y} \left\{ \cos\{\omega[t - n(\omega)x/c]\} + \cos\{\omega'[t - n(\omega')x/c]\} \right\}$   
 $= 2E_0 \hat{y} \cos\left\{\frac{1}{2}(\omega + \omega')t - \frac{1}{2} \frac{\omega n(\omega) + \omega' n(\omega')}{c} x\right\} \cos\left\{\frac{1}{2}(\omega' - \omega)t - \frac{\omega' n(\omega') - \omega n(\omega)}{2c} x\right\}$

$$\simeq 2E_0 \hat{y} \cos\{\omega_c [t - n(\omega_c)x/c]\} \cos\left\{\frac{1}{2} \Delta \omega [t - \frac{\omega' n(\omega') - \omega n(\omega)}{\omega' - \omega} x/c]\right\}$$

Carrier:

$$\text{Phase Velocity} = \frac{c}{n(\omega_c)}$$

envelope:

$$\text{Group Velocity} = \frac{c}{\frac{d}{d\omega} [\omega n(\omega)]}$$

Of course,  $\frac{d}{d\omega} [\omega n(\omega)] = n(\omega_c) + \omega_c n'(\omega_c)$ , when the derivative is evaluated at the center frequency  $\omega_c$ .